

大阪大学
社会人教育プログラム

FLAPW法の応用

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OUTLINE

- **Precision of the FLAPW Method**
 - **Wave Functions**
 - **Electron Density and Potential**
 - **Perturbative Consideration**
 - **Choice of MT Sphere Radius**
- **Typical Applications**
- **Summary**

WAVE FUNCTIONS

- **LAPW Basis**

$$\tilde{\phi}^{\mathbf{k}+\mathbf{K}}(\mathbf{r}) = \frac{1}{\sqrt{\Omega}} e^{i(\mathbf{k}+\mathbf{K})\cdot\mathbf{r}}$$

$$\phi^{\mathbf{k}+\mathbf{K}}(\mathbf{r}) = \sum_{\alpha lm} \left[A_{\alpha lm}^{\mathbf{k}+\mathbf{K}} R_l(r_\alpha) + B_{\alpha lm}^{\mathbf{k}+\mathbf{K}} \dot{R}_l(r_\alpha) \right] i^l Y_{lm}(\hat{\mathbf{r}}_\alpha)$$

- **Degrees of Variational Freedom**

$$K_{\max} \qquad l_{\max}$$

- **Choice of MT Sphere Radius**

ELECTRON DENSITY & POTENTIAL

$$\tilde{n}(\mathbf{r}) = \sum_{\mathbf{G}} e^{i\mathbf{G}\cdot\mathbf{r}} n_{\mathbf{G}} \quad n(\mathbf{r}) = \sum_{\alpha LM} n_{\alpha LM}(r_{\alpha}) i^L Y_{LM}(\hat{\mathbf{r}}_{\alpha})$$

$$\tilde{v}(\mathbf{r}) = \sum_{\mathbf{G}} e^{i\mathbf{G}\cdot\mathbf{r}} v_{\mathbf{G}} \quad v(\mathbf{r}) = \sum_{\alpha LM} v_{\alpha LM}(r_{\alpha}) i^L Y_{LM}(\hat{\mathbf{r}}_{\alpha})$$

- **Accuracy of Expansion**

$$G_{\max} \quad L_{\max}$$

variational parameters?

- **Choice of MT Sphere Radius**

PERTURBATIVE CONSIDERATION

- **Second-Order Perturbation**

$$\Delta_{\varepsilon^{\mathbf{k}}} = - \frac{|\langle \mathbf{k} + \mathbf{K} | \mathcal{H} | \mathbf{k} \rangle|^2}{|\mathbf{k} + \mathbf{K}|^2 - \varepsilon^{\mathbf{k}}}$$

- **Variational Parameters of the Wave Functions**

$$|\mathbf{k} + \mathbf{K}| \leq K_{\max}$$

$$l_{\max}$$

MT SPHERE RADIUS

When a sufficient $l_{\max}$ is assumed,

- In case of large MT sphere radius, because of smaller volume in the interstitial region fewer PW expansion is needed.
- In case of small MT sphere radius, because of larger volume in the interstitial region more PW expansion is needed.
- **A variational dimensionless parameter**

$$RK_{\max}$$

WAVE FUNCTIONS & ELECTRON DENSITY

$$n(\mathbf{r}) = \sum_{\mathbf{k}, n} |\psi_n^{\mathbf{k}}(\mathbf{r})|^2$$

$$G_{\max} \geq 2K_{\max}$$

$$L_{\max} \geq 2l_{\max}$$

- **Convergency of the electron density and potential expansion should be checked, especially when GGA is used.**
- **For small MT spheres used, higher $G_{\max}$ may be required to represent pseudized charge density.**

MT SPHERE RADIUS

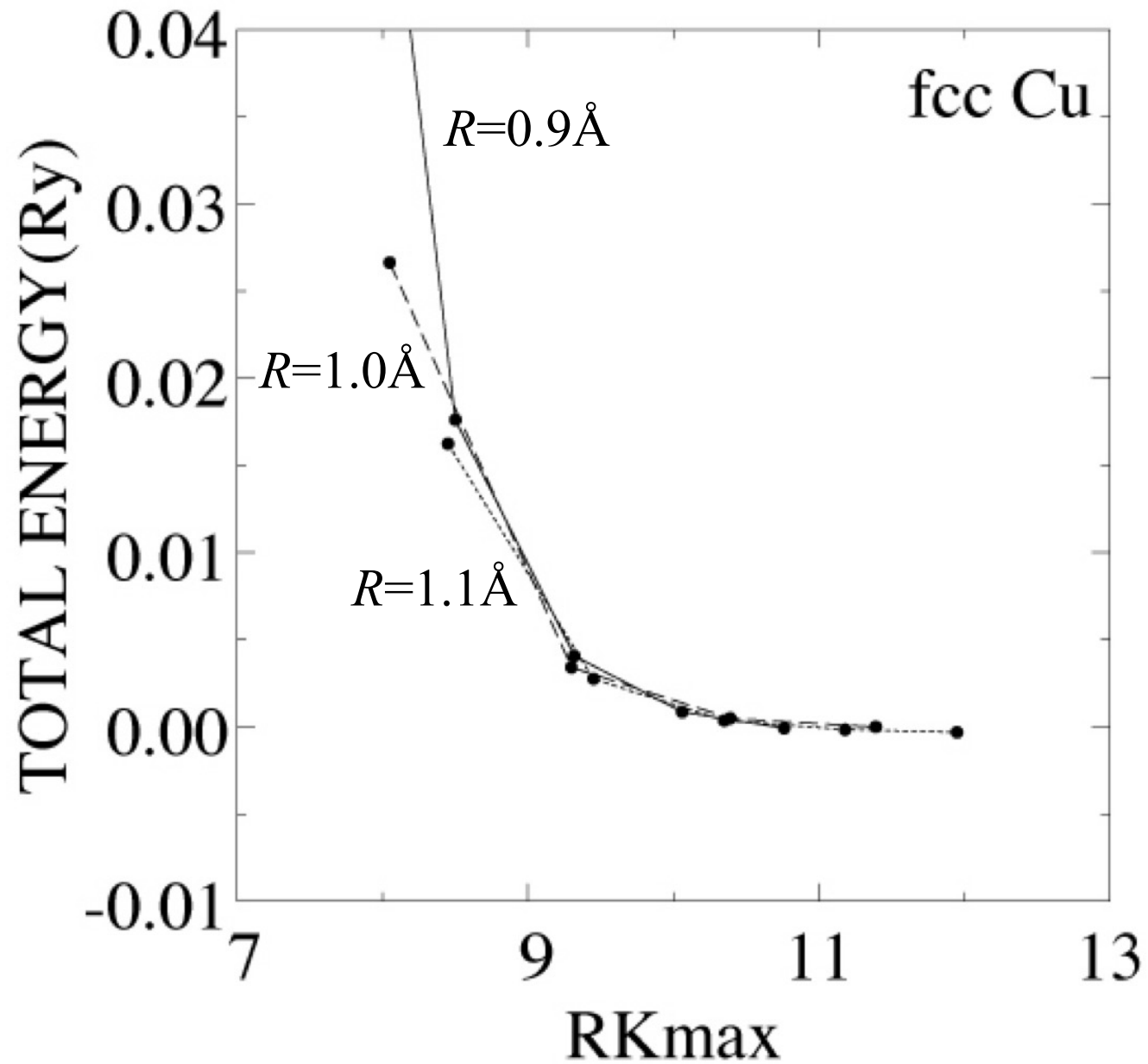
- **Non-overlapping spheres**

A margin should be considered if the atomic positions are changed, for example in a structural optimization calculation.

- **Negligible penetration of the core functions outside**

Within both assumptions with sufficient $l_{\max}$ and $L_{\max}$, the accuracy does not depend on the choice of MT sphere radius but does on $RK_{\max}$.

TOTAL ENERGY vs. MT RADIUS



PRECISION OF FLAPW METHOD

- **Wave Functions** $RK_{\max}$

- **Electron Density and Potential**

$$G_{\max} \geq 2K_{\max} \quad L_{\max} \geq 2l_{\max}$$

- **Choice of MT Sphere Radius**

- **Over-completeness of APW Basis Functions**

Since the PW basis is a complete set in all the space, an APW basis with the excessive number of PW results in indefinite solutions.

ELECTRIC FIELD GRADIENT

- **Definition of EFG**

$$V_{ij} = \frac{\partial^2 V(\mathbf{r})}{\partial i \partial j} \quad (i, j = x, y, z)$$

- **Maximum Term**

$$q = V_{ZZ}$$

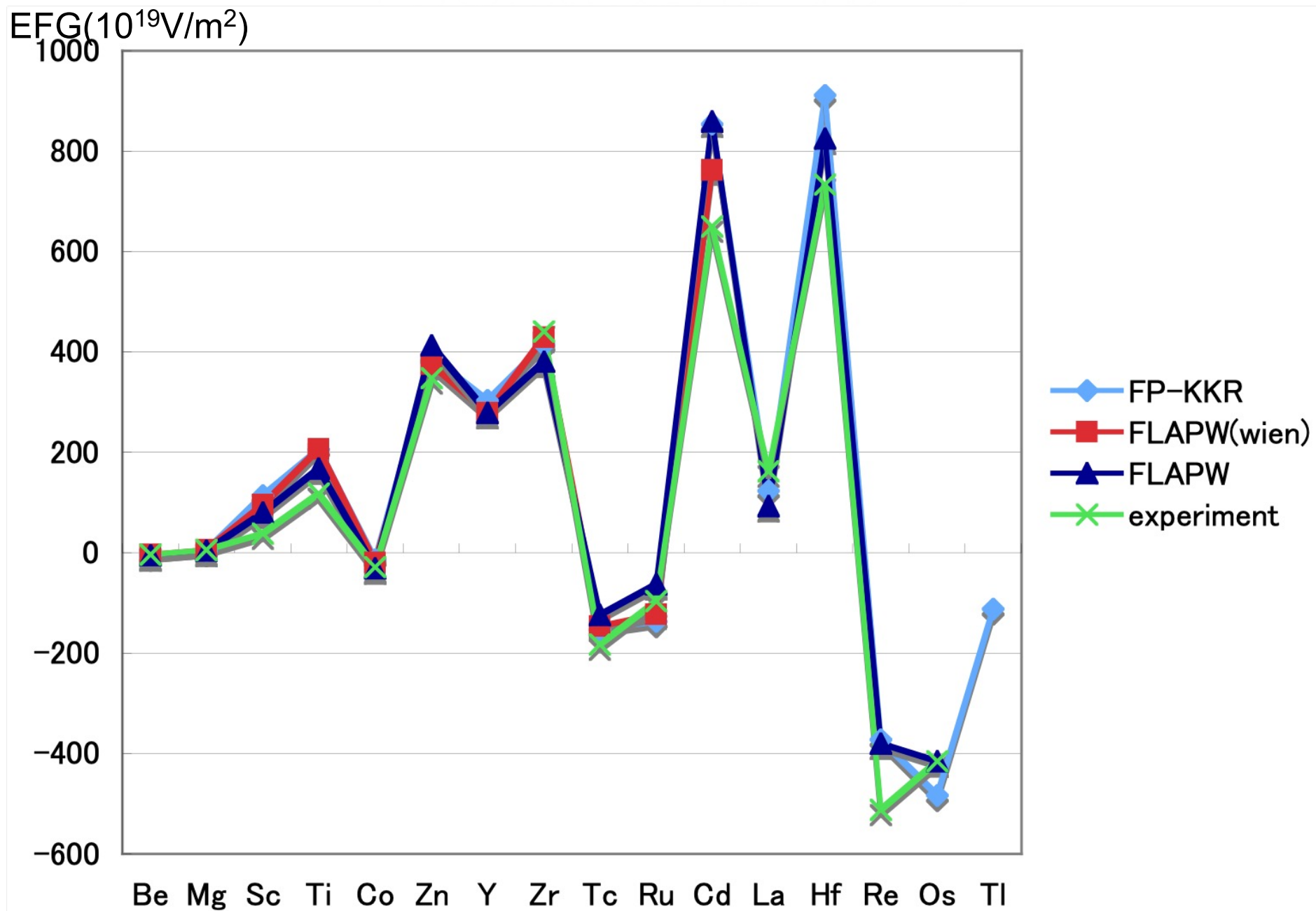
- **Asymmetric Parameter**

$$\eta = \frac{V_{XX} - V_{YY}}{V_{ZZ}}$$

- **Nuclear Magnetic Resonance Measurements**

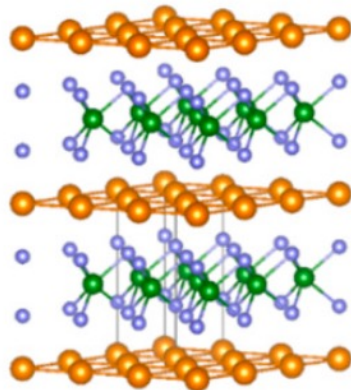
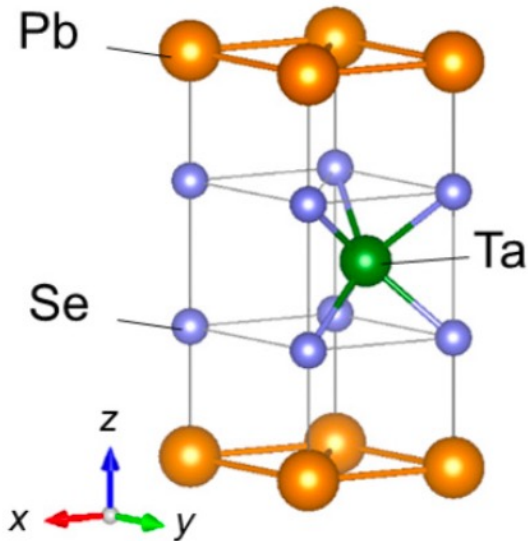
$$\nu_Q = \frac{3eqQ}{2I(2I-1)h} \quad \text{Quadrupole coupling frequency}$$

EFG IN HCP METALS

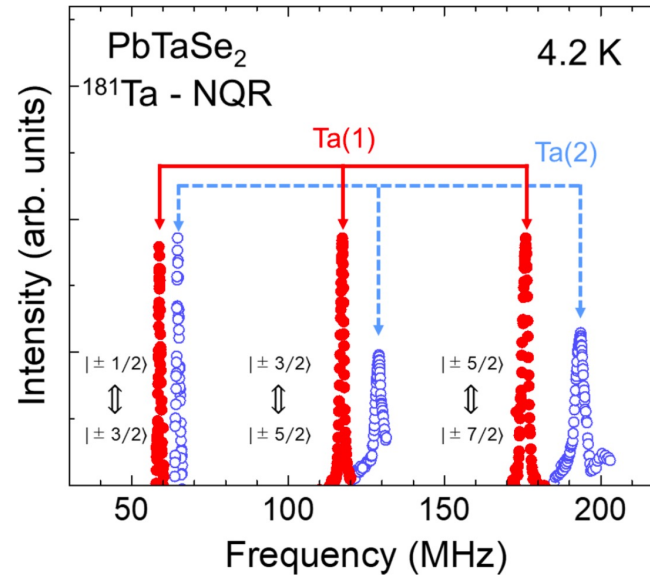


EFG at Ta in PbTaSe₂

PRB **102**, 214504 (2020)



P-6m2 structure



¹⁸¹Ta NQR spectra

	ν_Q (MHz)	η
Experiment		
Ta(1) site	58.7 ± 0.5	0.00–0.02
Ta(2) site	64.5 ± 0.5	0.03–0.04
DFT calculation		
$P\bar{6}m2$	57.78	0.00
$P6_3/mmc$	46.39	0.00

REFERENCE: EFG

- **EFG Calculations by FLAPW**
 - **PRL 54, 1192 (1985) Li_3N**
 - **PRB 37, 2792 (1988) hcp metals**
 - **PRB 102, 214504 (2020) PbTaSe_2**
- **EFG Calculations by FP-KKR**
 - **M. Ogura, PhD Thesis, Osaka University (2004)**
- **Review on EFG in NMR**
 - **Solid State Physics Vol. 5, pp. 321.**

PHONON CALCULATION

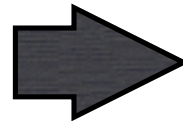
- **Equations of Motion** $M_\nu \ddot{u}_{n\nu\alpha} = - \sum_{n'\nu'\beta} \frac{\partial^2 E}{\partial u_{n\nu\alpha} \partial u_{n'\nu'\beta}} u_{n'\nu'\beta}$
- **Displacements** $u_{n\nu\alpha} = (M_\nu)^{-\frac{1}{2}} C_{\nu\alpha} e^{i(\mathbf{q} \cdot \mathbf{R}_n - \omega(\mathbf{q})t)}$

- **Dynamical Matrix**

$$D_{\nu\alpha,\nu'\beta}(\mathbf{q}) = (M_\nu M_{\nu'})^{-\frac{1}{2}} \sum_{n'} \frac{\partial^2 E}{\partial u_{0\nu\alpha} \partial u_{n'\nu'\beta}} e^{i\mathbf{q} \cdot \mathbf{R}_{n'}}$$

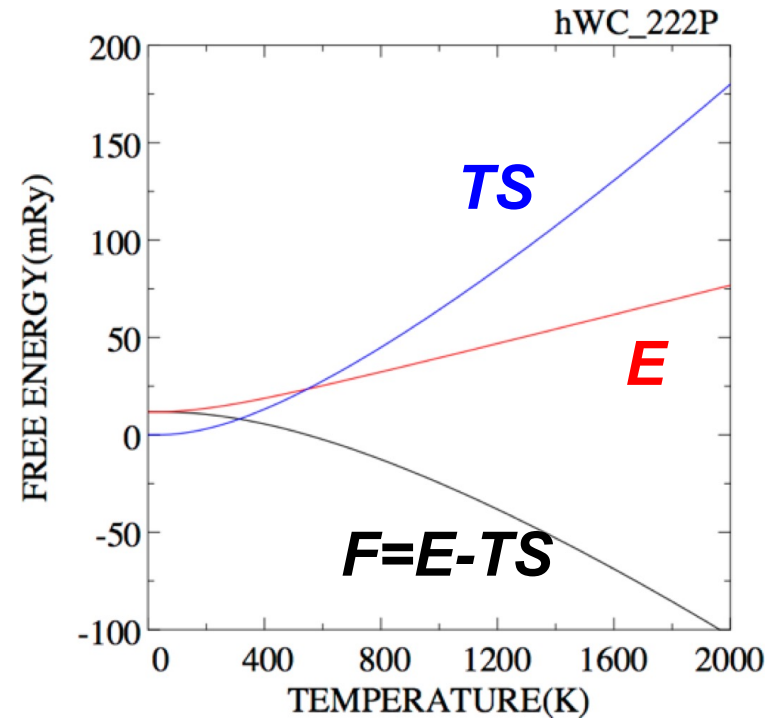
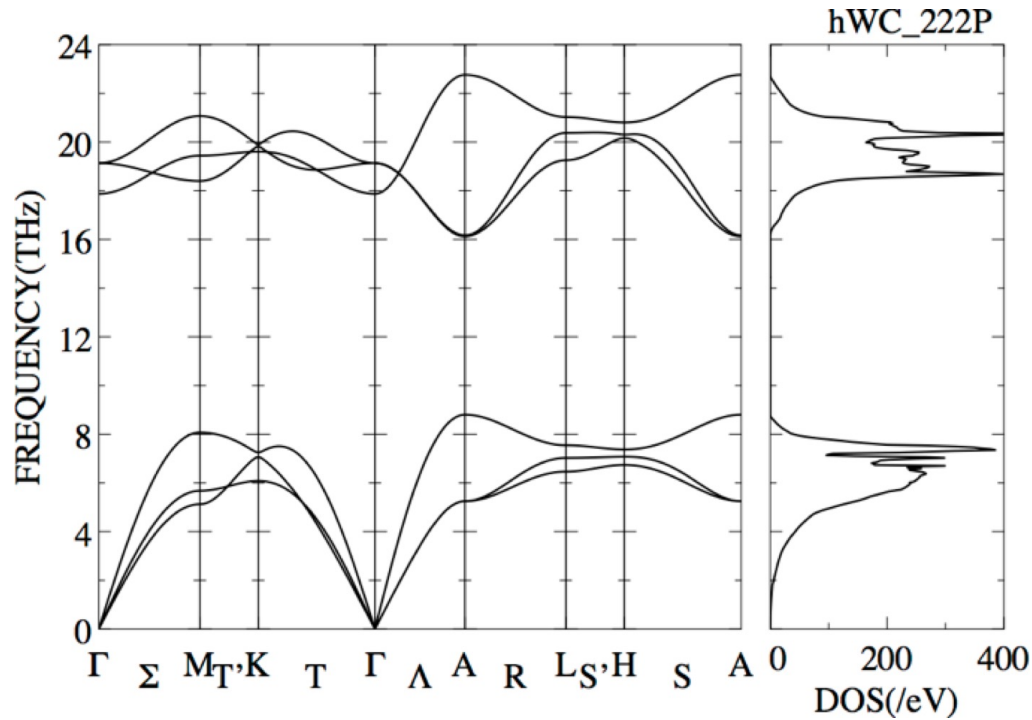
$$\approx -(M_\nu M_{\nu'})^{-\frac{1}{2}} \frac{F_{0\nu\alpha}}{u_{0\nu'\beta}} \left(\begin{array}{l} u_{n\nu\alpha} = u_{0\nu\alpha} e^{i\mathbf{q} \cdot \mathbf{R}_n} \\ F_{n\nu\alpha} = -\frac{\partial E}{\partial u_{n\nu\alpha}} \end{array} \right.$$

$$\mathbf{D}(\mathbf{q})\mathbf{C} = \mathbf{C}\omega^2(\mathbf{q})$$



Phonon modes

PHONON CALCULATION OF WC



$$E = \sum_{n,\mathbf{q}} \hbar\omega_n(\mathbf{q}) \left[\frac{1}{2} + (e^{\beta\hbar\omega_n(\mathbf{q})} - 1)^{-1} \right]$$

$$F = \sum_{n,\mathbf{q}} \left[\frac{\hbar\omega_n(\mathbf{q})}{2} + \beta^{-1} \ln \left(1 - e^{-\beta\hbar\omega_n(\mathbf{q})} \right) \right]$$

ELASTIC AND DYNAMICAL STABILITY OF QUATERNARY HEUSLER

CoCrScAl

Born-Cauchy conditions

➤ **Elastic stability**

Bulk modulus : $K = (C_{11} + 2C_{12})/2$

Shear modulus : $C_s = (C_{11} - C_{12})/2$

C_{44}

(GPa)

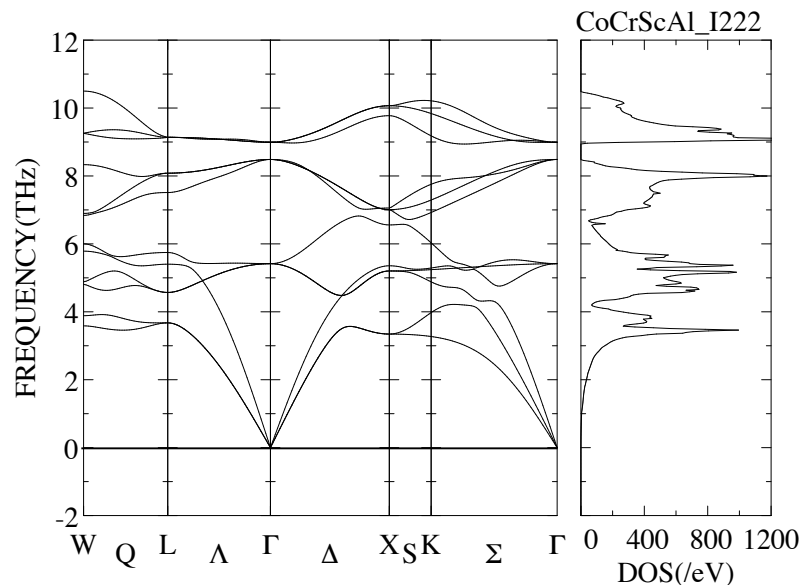
116.1 > 0

61.9 > 0

103 > 0

Elastically stable

➤ **Dynamical stability**



No soft mode



Dynamically stable

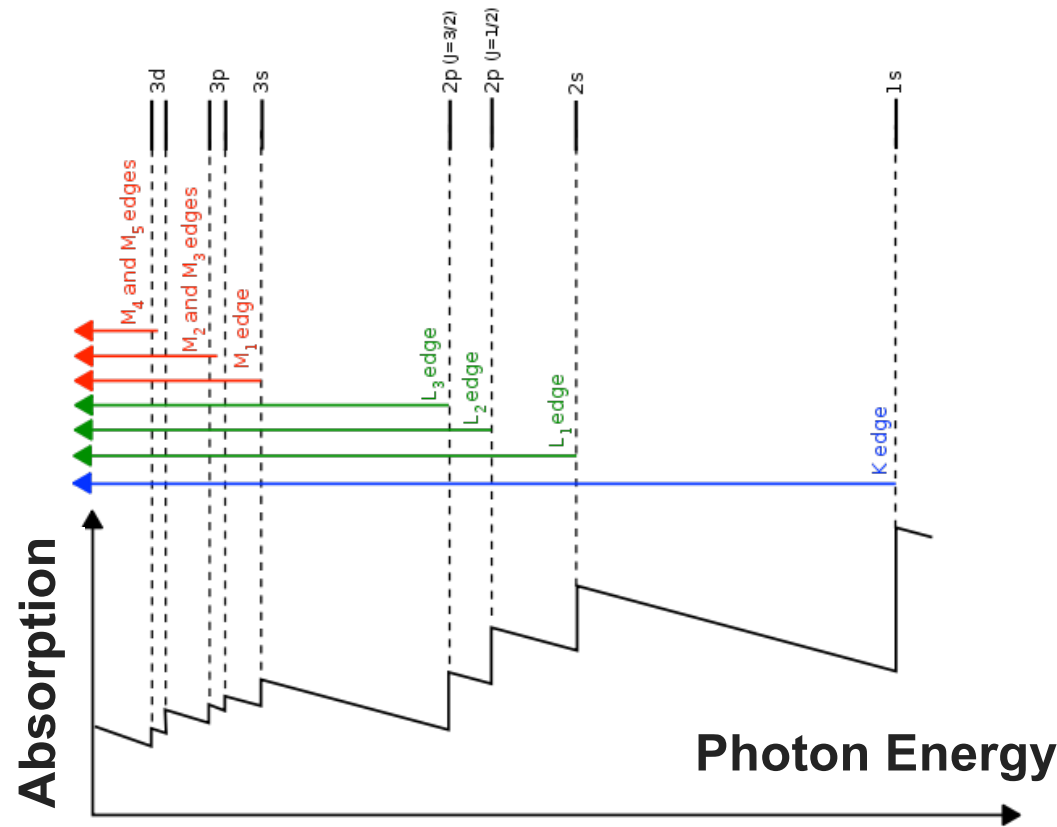
**F. Kuroda, T. Fukushima, TO,
J. Appl. Phys. 127, 193904 (2020).**

REFERENCE: ATOMIC FORCES BY FLAPW

- **Soler, Williams, PRB40, 1560 (1989); PRB42, 9728 (1990); PRB47, 6784 (1993).**
- **Yu, Singh, Krakauer, PRB43, 6411 (1991); PRB45, 8671 (1992).**
- **Goedecker, Maschke, PRB45, 1597 (1992).**
- **T. Oguchi, “Interatomic Potentials and Structural Stability”, ed. by K. Terakura and H. Akai, (Springer-Verlag, 1993) p. 33.**
- **D. J. Singh, “Planewaves, pseudopotentials and the LAPW method” (Kluwer Academic, Boston, 1994).**

X-ray Absorption Spectroscopy (XAS)

- X-ray absorption: electronic transitions by photon from inner core states to empty states
 - suitable for studying valence states with specific **element**, **orbital** and **symmetry** by tuning the photon **energy** and **polarization**



Fundamentals of XAS

- **Transition Probability by Fermi Golden Rule**

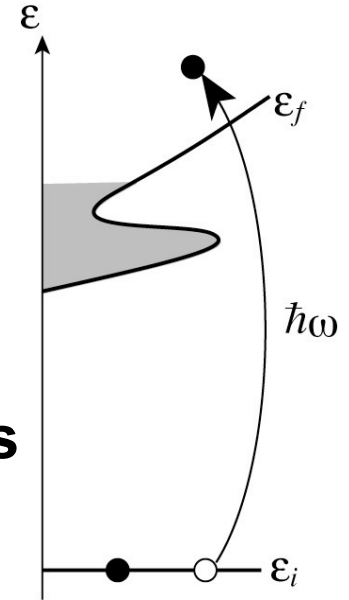
$$W(\omega) = \frac{2\pi}{\hbar} \sum_{f,i} |\langle f | \hat{F} | i \rangle|^2 \delta(\varepsilon_f - \varepsilon_i - \hbar\omega)$$

- **Rigorous Theory**

ε_i ε_f **Total energies of initial and final states**

$|i\rangle$ $|f\rangle$ **Many-body initial and final states**

$\hat{F}$ **Perturbation by photon in many-body hamiltonian**



- **Quasi-particle (One-electron) Approach**

ε_i ε_f **Orbital energies of initial and final states**

$|i\rangle$ $|f\rangle$ **One-electron initial and final states**

$\hat{F}$ **Perturbation by photon in one-electron hamiltonian**

Theoretical Approaches

- **Quasi-particle Approaches**

- **Ground-state calculation**

- **Core-hole calculation**

- **Perturbation calculation (GW, ...)**

- **.....**

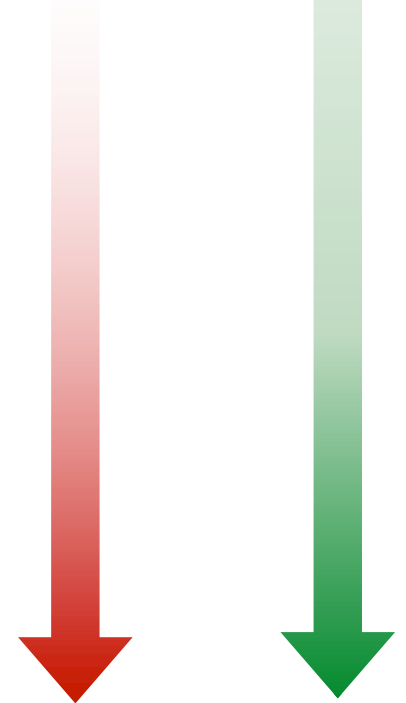
- **Δ SCF Approach**

- **(Excited State) – (Ground State)**

- **Many Body Approaches**

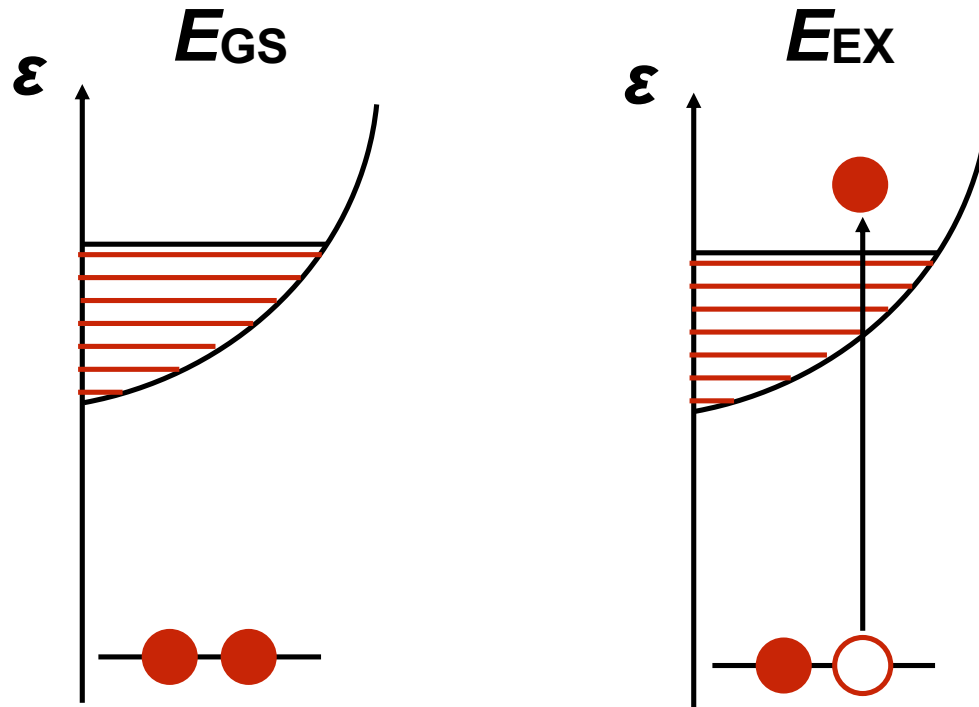
**spectrum shape
polarization dependence**

excitation energy



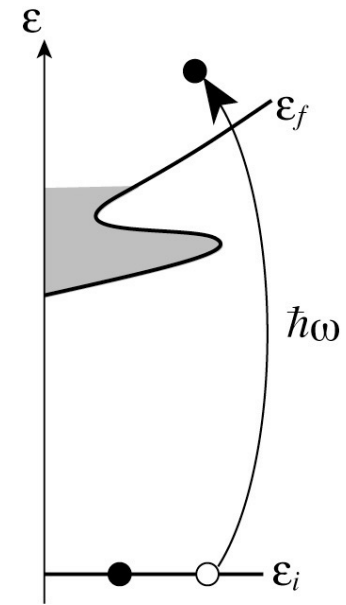
Δ SCF

- Self-Consistent-Field (SCF) DFT Calculation
 - Grand state: E_{GS}
 - Excited state: E_{EX}
- Excitation Energy = $E_{EX} - E_{GS}$



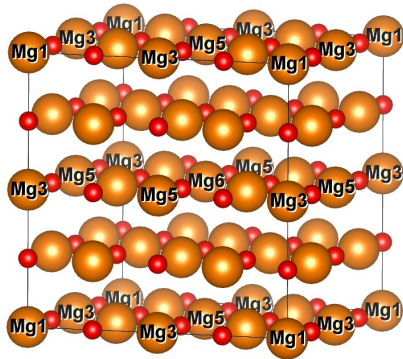
Supercell Model

- Core hole: well localized
- Excited electron: usually delocalized
 - Hybridization with neighbors
 - Electrostatic fields due to holes and excited electrons at neighboring sites



hole concentration

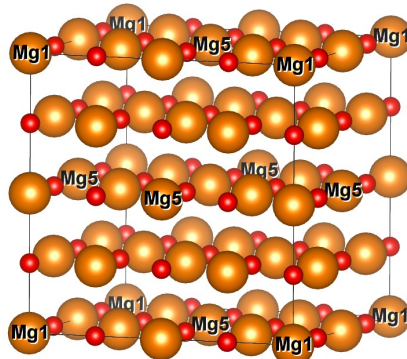
1/4



(MgO)₄

4³=64

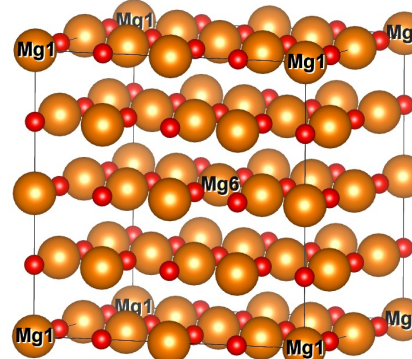
1/8



(MgO)₈

8³=512

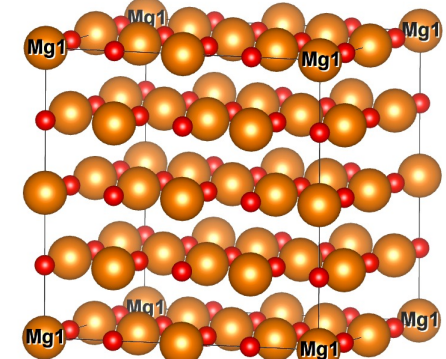
1/16



(MgO)₁₆

16³=4096

1/32



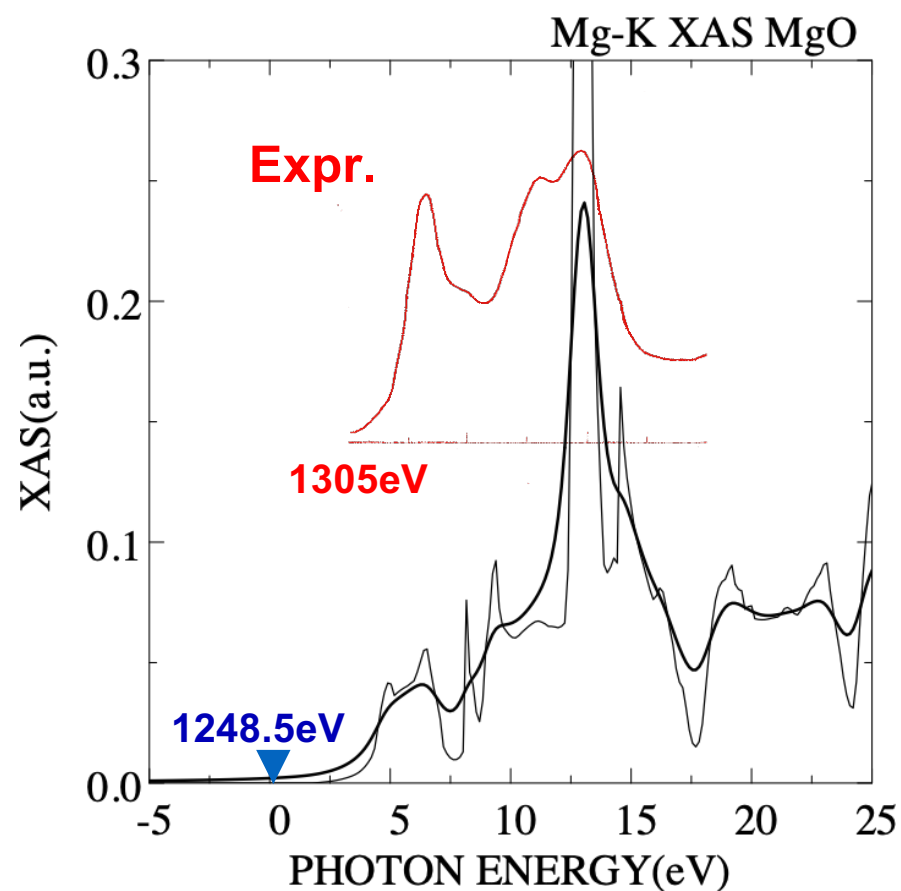
(MgO)₃₂

32³=32768

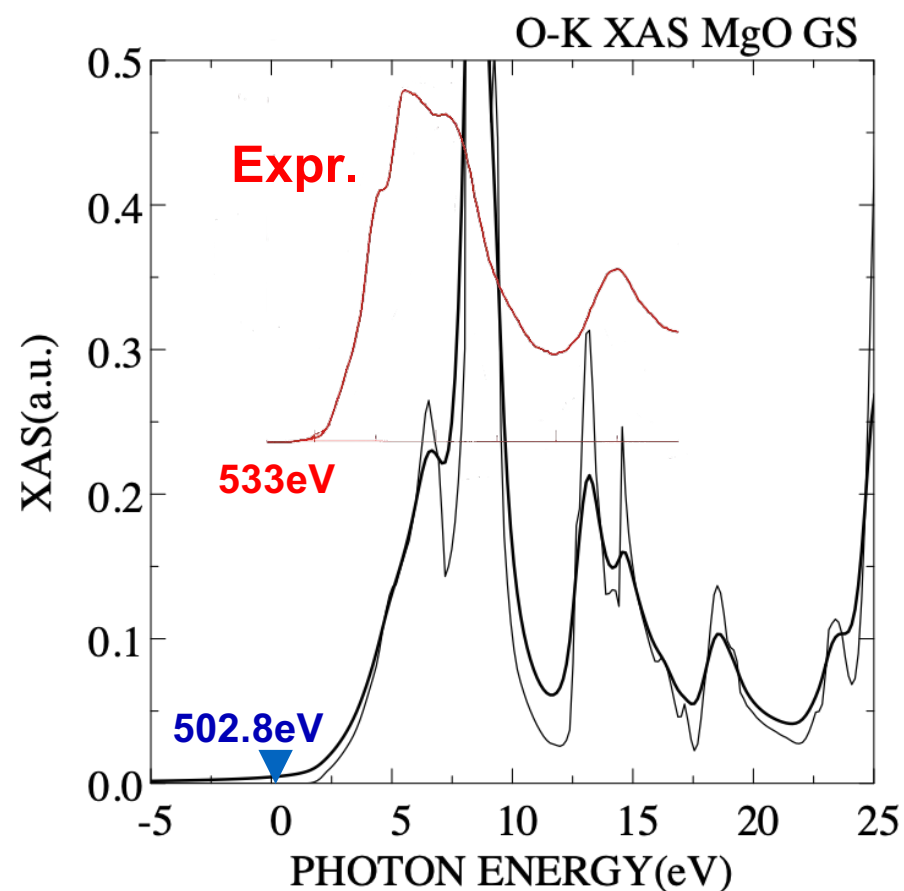
XAS of MgO

Grand-State Approximation

Mg K-edge (E1: $1s \rightarrow 3p$)

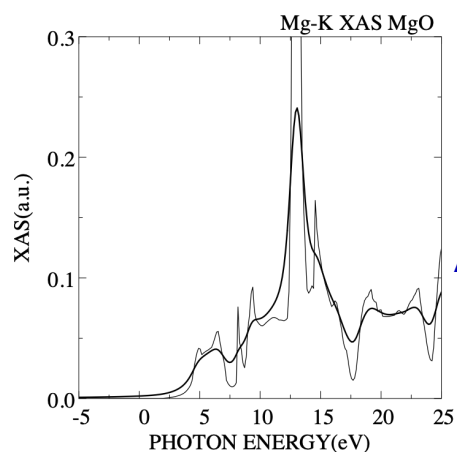


O K-edge (E1: $1s \rightarrow 2p$)

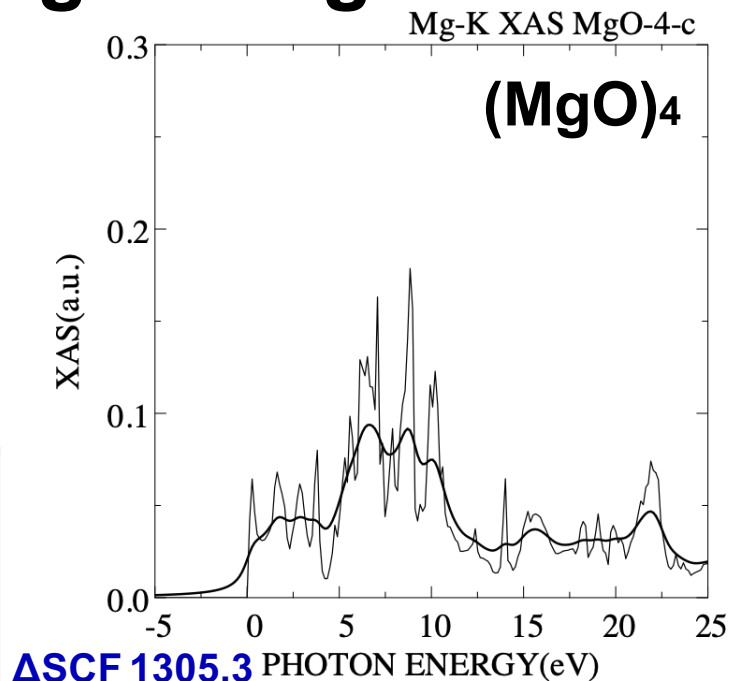


Thick lines: broadened by Lorentzian

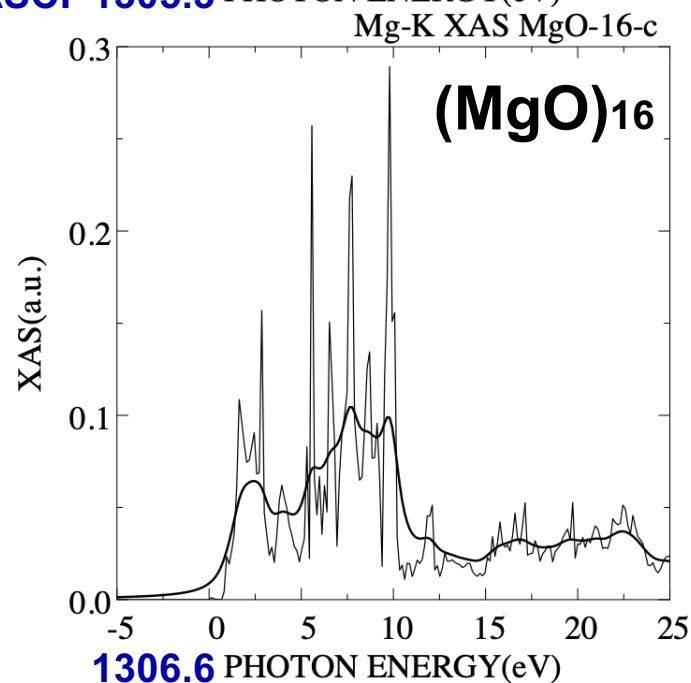
Mg K-Edge of MgO with **Core Hole** in Mg-1s



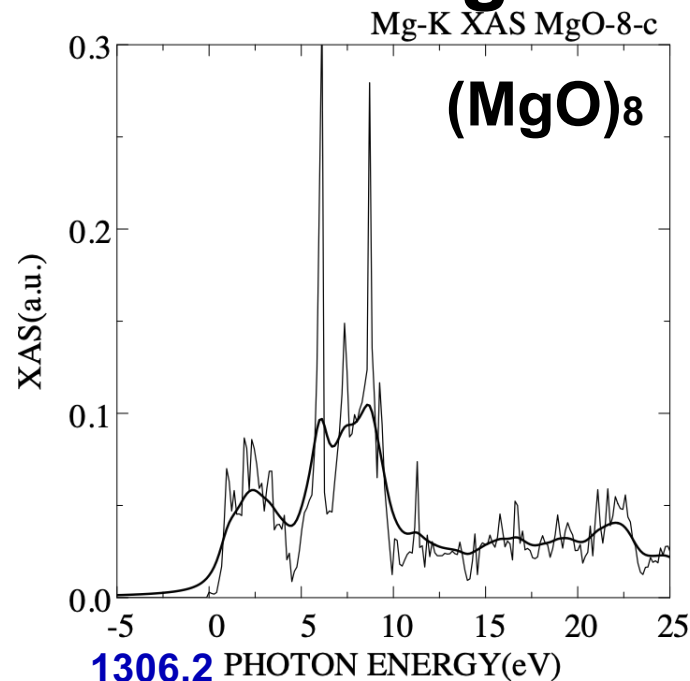
G.S.



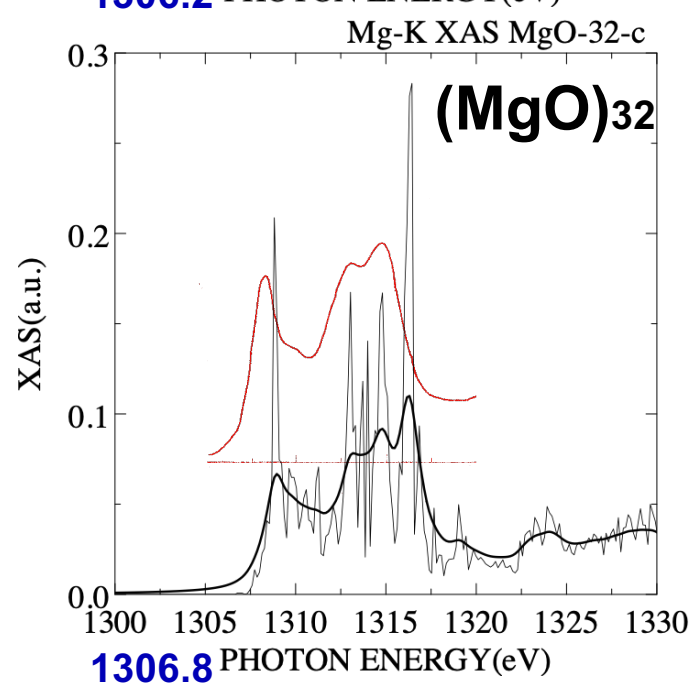
ΔSCF 1305.3



1306.6

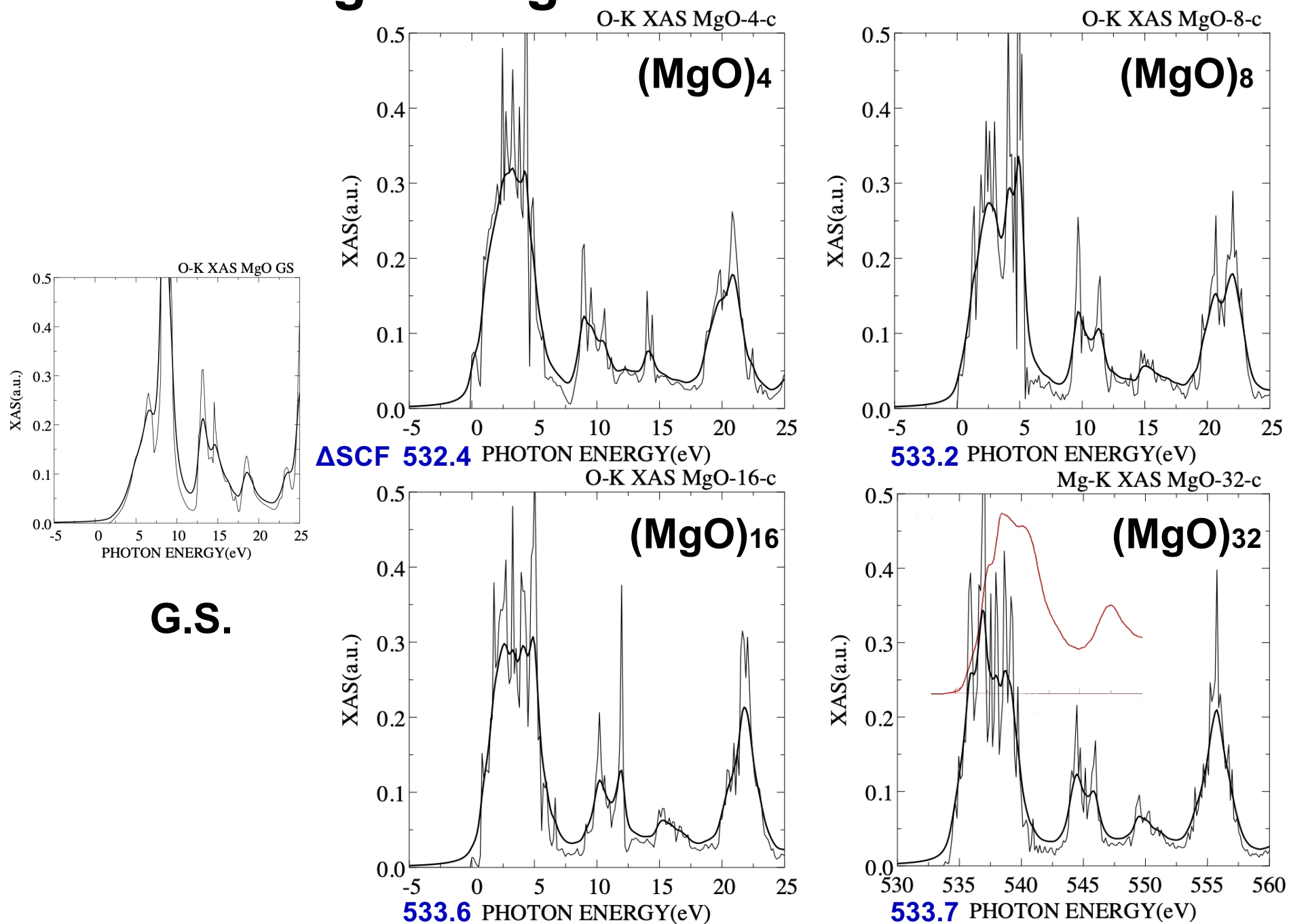


1306.2



1306.8

O K-Edge of MgO with **Core Hole** in O-1s



Summary of XAS

- **XAS calculations with core hole**
 - **Better shape of the spectra**
 - **Better excitation energy in the Δ SCF sense**
- **Test calculations at *K* edges of MgO**
- **Further test**
 - **Different edges (L, M, N, ...)**
 - **Different materials (magnetic, ...)**
 - **Effects on polarization (MCD, LD, ...)**

SUMMARY

- **FLAPW Method**
 - **Precise and robust, but often heavy**
 - **$\text{CPU} \sim N_{\text{atom}}^3$**
 - **Complementary role with PP-PW**
 - **Care of the convergency in several parameters involved**
- **A wide variety of applications of FLAPW**

First-Principles Calculations

